

The quantum $ISL(2,C)$ group

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

1994 J. Phys. A: Math. Gen. 27 7099

(<http://iopscience.iop.org/0305-4470/27/21/024>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 171.66.16.68

The article was downloaded on 01/06/2010 at 23:08

Please note that [terms and conditions apply](#).

The quantum $ISL(2, \mathbb{C})$ group

Paweł Maślanka†‡

Department of Functional Analysis, Institute of Mathematics, University of Łódź,
ul. St. Banacha 22, 90-238 Łódź, Poland

Received 29 March 1994

Abstract. The classical r -matrix implied by the quantum κ -Poincaré algebra of Lukierski, Nowicki and Ruegg is used to generate a Poisson structure on the $ISL(2, \mathbb{C})$ group. A quantum deformation of the $ISL(2, \mathbb{C})$ group (on the Hopf algebra level) is obtained by a trivial quantization.

Recently a quantum deformation of Poincaré algebra has been constructed which depends on dimensional parameter κ [1]. It has attractive properties like, for example, space isotropy or the existence of a natural cut-off κ . Some physical consequences of deformed Poincaré symmetry were considered by Bacry [2]. The global counterpart of this algebra, i.e. the quantum Poincaré group, was constructed by Zakrzewski [3]. However, as in the ‘classical’ case, for many purposes (the unitary representations of half-integer spin, supersymmetric extensions, etc) it is important to consider the universal enveloping group $ISL(2, \mathbb{C})$. It is the aim of the present paper to give a description of this group. To this end we use the approach of Zakrzewski and introduce the co-Poisson structure determined from $(1/\kappa)$ -expansion of the antisymmetric part of the coproduct of the κ -Poincaré algebra. By duality the Poisson structure on $ISL(2, \mathbb{C})$ is obtained and quantized. Due to the form of Poisson brackets no quantization ambiguities appear. Therefore one can expect that the ‘quasiclassical’ (in $1/\kappa$) duality can be extended to the full quantum duality (this is the case for $E_q(2)$ (see [4])).

The classical $ISL(2, \mathbb{C})$ group consists of the pairs (a, A) , where a is a fourvector and A is a matrix of the $SL(2, \mathbb{C})$ group and with the composition law

$$(a, A) * (a_1, A_1) = (a + \Lambda(A)a_1, AA_1) \quad (1)$$

where

$$\Lambda(A) = [\Lambda_\nu^\mu(A)]_{\mu, \nu=0}^3 = \left[\frac{1}{2} \text{Tr}(\sigma_\mu A \sigma_\nu A^+) \right]_{\mu, \nu=0}^3 \quad (2)$$

is the Lorentz matrix corresponding to the matrix A (and σ_μ are the Pauli matrices).

In order to determine the Poisson structure on the $ISL(2, \mathbb{C})$ group we have to calculate the right- and left-invariant vector fields. To do this we can consider first the $GL(2, \mathbb{C})$ group. Due to the fact that the r -matrix, which determines the Poisson bracket by the formula

$$\{f, g\} = 2r^{\alpha\beta} (X_\alpha^R f X_\beta^R g - X_\alpha^L f X_\beta^L g) \quad (3)$$

† E-mail address: PMAŚLAN@PLUNLO51.BITNET

‡ Supported by KBN grant 2 0218 91 01.

where X_α^R, X_β^L are the right- and left-invariant vector fields, contains only the generators of $SL(2, \mathbb{C})$ we can consider all matrix elements A_β^α as independent and impose the unimodularity condition at the very end (i.e. $\det A$ has vanishing Poisson bracket with everything). In this way we obtain the following expressions for the invariant vector fields:

$$\begin{aligned}
 (X^L)_\mu &= \Lambda_\mu^\nu(A) \frac{\partial}{\partial a^\nu} \\
 (X^L)_\gamma^\alpha &= A_\gamma^\beta \frac{\partial}{\partial A_\alpha^\beta} \\
 (\bar{X}^L)_\gamma^\alpha &= \bar{A}_\gamma^\beta \frac{\partial}{\partial \bar{A}_\alpha^\beta} \\
 (X^R)_\mu &= \frac{\partial}{\partial a^\mu} \\
 (X^R)_\beta^\alpha &= A_\gamma^\alpha \frac{\partial}{\partial A_\gamma^\beta} + \frac{1}{2}(\sigma_\mu \sigma_\nu)_\beta^\alpha a^\mu \frac{\partial}{\partial a^\nu} \\
 (X^R)_\beta^\alpha &= \bar{A}_\gamma^\alpha \frac{\partial}{\partial \bar{A}_\gamma^\beta} + \frac{1}{2}(\bar{\sigma}_\mu \bar{\sigma}_\nu)_\beta^\alpha a^\mu \frac{\partial}{\partial a^\nu}
 \end{aligned}
 \tag{4}$$

$(\mu, \nu = 0, 1, 2, 3; \alpha, \beta = 1, 2)$.

In the Lie algebra basis corresponding to the above vector fields the r -matrix obtained by Zakrzewski [3] takes the form

$$r = -\frac{i}{\kappa}(L_k \otimes P_k - P_k \otimes L_k)
 \tag{5}$$

where

$$\begin{aligned}
 L_k &= \frac{1}{2i}[(\sigma_k)_\beta^\alpha X_\alpha^\beta + (\bar{\sigma}_k)_\beta^\alpha \bar{X}_\alpha^\beta] \\
 P_k &= X_k
 \end{aligned}$$

and $k = 1, 2, 3$.

This r -matrix induces (by duality) the following Poisson brackets of the functions on $ISL(2, \mathbb{C})$:

$$\begin{aligned}
 \{f, g\} &= -\frac{1}{2\kappa} \{[(\sigma_k)_\beta^\alpha (X^R)_\alpha^\beta f + (\bar{\sigma}_k)_\beta^\alpha (\bar{X}^R)_\alpha^\beta f](X^R)_k g \\
 &\quad - [(\sigma_k)_\beta^\alpha (X^R)_\alpha^\beta g + (\bar{\sigma}_k)_\beta^\alpha (\bar{X}^R)_\alpha^\beta g](X^R)_k f \\
 &\quad - [(\sigma_k)_\beta^\alpha (X^L)_\alpha^\beta f + (\bar{\sigma}_k)_\beta^\alpha (\bar{X}^L)_\alpha^\beta f](X^L)_k g \\
 &\quad + [(\sigma_k)_\beta^\alpha (X^L)_\alpha^\beta g + (\bar{\sigma}_k)_\beta^\alpha (\bar{X}^L)_\alpha^\beta g](X^L)_k f\}.
 \end{aligned}
 \tag{6}$$

If we perform the standard quantizations of the Poisson brackets of the coordinate functions on $ISL(2, \mathbb{C})$, by replacing $\{, \}$ $\rightarrow \frac{1}{i}[,]$, one obtains the following set of

commutation relations

$$\begin{aligned}
 [A_\beta^\alpha, A_\delta^\gamma] &= 0 & [\bar{A}_\beta^\alpha, \bar{A}_\delta^\gamma] &= 0 & [\bar{A}_\beta^\alpha, A_\delta^\gamma] &= 0 \\
 [a^k, a^j] &= 0 & [a^k, a^0] &= \frac{i}{\kappa} a^k \\
 [A_\beta^\alpha, a^0] &= \frac{i}{2\kappa} (A\sigma_k)_\beta^\alpha \Lambda_k^0(A) \\
 [A_\beta^\alpha, a^k] &= \frac{i}{2\kappa} [(A\sigma_n)_\beta^\alpha \Lambda_n^k(A) - (\sigma_k A)_\beta^\alpha] \\
 [\bar{A}_\beta^\alpha, a^0] &= \frac{i}{2\kappa} (\bar{A}\bar{\sigma}_k)_\beta^\alpha \Lambda_k^0(A) \\
 [\bar{A}_\beta^\alpha, a^k] &= \frac{i}{2\kappa} [(\bar{A}\bar{\sigma}_n)_\beta^\alpha \Lambda_n^k(A) - (\bar{\sigma}_k \bar{A})_\beta^\alpha].
 \end{aligned} \tag{7}$$

Since the composition law is compatible with Poisson bracket, the above commutation rules are compatible with the following coproduct:

$$\begin{aligned}
 \Delta(A_\beta^\alpha) &= A_\gamma^\alpha \otimes A_\beta^\gamma \\
 \Delta(a^\mu) &= \Lambda_\nu^\mu(A) \otimes a^\nu + a^\mu \otimes I.
 \end{aligned} \tag{8}$$

The antipode is given by

$$S((a, A)) = (-\Lambda_\nu^\mu(A^{-1})a^\nu, A^{-1}). \tag{9}$$

We conclude that relations (7)–(9) define a Hopf algebra, the $ISL_\kappa(2, \mathbb{C})$ group. Note that the map

$$(a, A) \rightarrow (a, \Lambda(A)) \tag{10}$$

is the Hopf algebra homomorphism from $ISL_\kappa(2, \mathbb{C})$ to the κ -Poincaré group of Zakrzewski.

Acknowledgment

I am grateful to Professor P Kosiński for discussions.

References

- [1] Lukierski J, Nowicki A, Ruegg H and Tolstoy V 1991 *Phys. Lett.* **264B** 331
 Güller S, Kosiński P, Kunz J, Majewski M and Maślanka P 1992 *Phys. Lett.* **286B** 57
 Lukierski J, Nowicki A and Ruegg H 1992 *Phys. Lett.* **344B** 344
- [2] Bacry H 1993 *Phys. Lett.* **306B** 44; 1992 Classical electrodynamics on a quantum Poincaré group *Preprint CPT-92/P.2837*; 1993 Which deformation of the Poincaré group? *CPT Preprint*; 1993 Do we have to believe in dark matter? *CPT Preprint*
- [3] Zakrzewski S 1993 Quantum Poincaré group related to k -Poincaré algebra *Warsaw University Preprint J. Phys. A: Math. Gen.* **27** 2075)
- [4] Maślanka P 1994 The two dimensional quantum euclidean group 2/93 *IMUL Preprint*; 1994 *J. Math. Phys.* **35** 1976
 Ballesteros A, Celeghini E, Giachetti R, Sorace E and Tarlini M 1994 An R -matrix approach to the quantization of the euclidean group $E(2)$ *J. Phys. A: Math. Gen.* **26** 7495